

EXPERIMENTAL WALL CORRECTION FACTORS OF SINGLE SOLID SPHERES IN TRIANGULAR AND SQUARE CYLINDERS, AND PARALLEL PLATES

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(Received 28 August 1979; in revised form 27 June 1980)

Abstract—To examine wall effects, the steady state settling velocities of single solid spheres of various diameters were measured in triangular and square cylinders, and between parallel plates, which were filled up with highly viscous aqueous millet jelly solutions. The wall effect for the steady state motion of single spheres in a narrow space has to date been correlated by the wall correction factor. The wall correction factors of single solid spheres were correlated in such apparatus in the creeping flow region.

INTRODUCTION

Particles interact with other particles in hindered settling and particulate fluidization. The interaction among the particles was correlated by the concept of the void function proposed by Ito (1957, 1962). The previous theoretical and experimental void functions were summarized by Ishii (1965). On the other hand, particles moving in a narrow space are influenced by the wall. The wall effect has been correlated by the concept of the wall correction factor.

There have been many theoretical and experimental studies of wall correction factors of single solid spheres in circular cylinders in the creeping flow region, which have been summarized by Iwaoka & Ishii (1979). Recently, Iwaoka & Ishii (1979) reported the wall correction factors of single solid spheres in circular cylinders in the creeping flow region and their data agreed very well with the exact solutions of Haberman (1956), and Paine & Scherr (1975).

However, there have been few experimental and theoretical studies on the wall correction factors of single solid spheres in various apparatus other than circular cylinders.

In this study, the wall correction factors of single solid spheres were measured in triangular and square cylinders, between parallel plates in the creeping flow region.

The purpose of this series of studies is to measure the wall effect of single solid spheres under steady state motion in various apparatus in the creeping flow region. These studies also resulted in development of new falling ball viscometers.

2. EXPERIMENTAL

The experimental apparatus used were made from transparent acryl resin plates and are shown in figures 1-3. The triangular cylinder with the triangular jacket was 100 cm in height and 5 mm in perpendicular distance from the center to each inner side of a regular triangular cross-section of the cylinder. The square cylinder with the square jacket was 100 cm in height and 1 cm in an inner side length. The parallel plates with the rectangular jacket were 100 cm in height, 0.99 cm in shorter width and 10 cm in longer width. The entrance and end regions were 20 cm long, based on previous work (Brenner 1961, Tanner 1963, Yu *et al.* 1979) resulting in a measuring section length of 60 cm. Unfortunately, curvature was observed in the lower sections of the triangular and square cylinder, and hence the measuring sections of these cylinders were reduced from 20 to 40 cm from the top.

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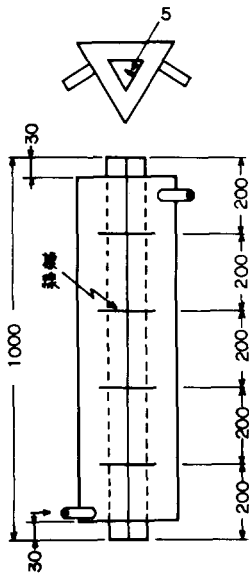


Figure 1. Triangular cylinder.

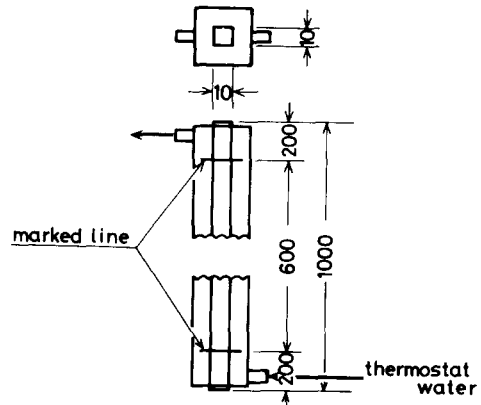


Figure 2. Square cylinder.

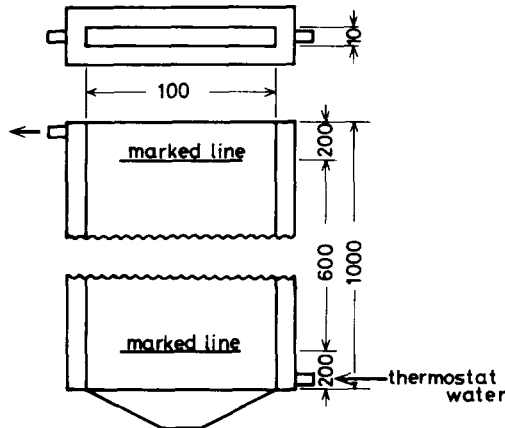


Figure 3. Parallel plates.

The apparatus were maintained precisely vertically by making it accurately parallel to a thread suspending a weight. The apparatus were filled with the aqueous millet jelly solution, which was highly viscous, but still a Newtonian fluid (Nakagawa 1975). For each run the aqueous millet jelly solutions in the apparatus were maintained at constant temperatures of 25 and 30°C by flowing water in thermostat into the jacket. The viscosities of the aqueous millet jelly solutions were measured at 20, 25, 30, 35 and 40°C by the Hoespler falling ball viscometer, and were 1805 and 809.0 centipoise at 25 and 30°C for the experiments with the triangular cylinder, respectively and were 4420 and 2490 centipoise at 25 and 30°C for the experiments with the square cylinder and parallel plates. The densities of the aqueous millet jelly solutions were measured at 20, 25, 30, 35 and 40°C by specific gravity bottle in a thermostat.

The single ball bearings of diameters of 1.00, 2.00, 3.00, 4.00, 5.00, 6.00, 7.00, 8.00 and 9.00 mm were dropped into the centers of the tops of the apparatus. The times for the single ball bearings to pass through the measuring sections of the apparatus were measured by a stopwatch. The experiments were repeated three times for each run. The ball bearings were conditioned in the weight measuring bottle filled with the same aqueous millet jelly solutions in the thermostat and at the same temperature as the experimental apparatus.

3. EXPERIMENTAL RESULTS AND DISCUSSIONS

The single ball bearings were observed to settle along the center line in the measuring sections of the apparatus with no rotation. In the case that the ball bearing was dropped along the inside wall at the top of the apparatus, the ball bearing approached the center line with rotation and settled with no rotation along the centerline of the apparatus within the entrance region. These phenomena are naturally reasonable, as the viscous dissipation, namely entropy production rate becomes the minimum in the case that the solid sphere settles along the centerline of the apparatus. The steady state settling velocities, u_w were calculated from the times for which the single ball bearings passed through the measuring sections of the apparatus.

In the creeping flow region the following relationship holds between the free settling velocity, u_F and the settling velocity influenced by the wall, u_w .

$$F(r) = \frac{u_F}{u_w} . \quad [1]$$

The free settling velocities were calculated from Stokes law, and the Reynolds numbers, of which representative velocities are free and hindered settling velocities, u_F and u_w , ranged from 0.000241 to 0.699 and from 0.000216 to 0.0395, respectively.

Figures 4 and 5 show semi-log plots of the wall correction factors of single solid spheres in triangular and square cylinders vs the dimensionless radii of the solid spheres, respectively.

Figure 6 shows a plot of the wall correction factors of single solid spheres in parallel plates vs the dimensionless radii of the solid spheres.

The clearance between the solid sphere of a fixed diameter and the wall of the apparatus increases with the order of the triangular and square cylinders, and parallel plates. The wall correction factors increases with decreasing clearance, as seen in figures 4–6. All the data negligibly scatter. In each run at least two data were measured and they negligibly scatter. Hence, one point of the data indicates at least two measurements in figures 4–6.

In the case of circular cylinders, the solid sphere of necessity settled with rotation within the very small clearances, even though precautions were taken. This rotation phenomena has been analysed by lubrication theory (Christopherson & Dowson 1959, Bungay & Brenner 1973). However, this rotation phenomena was rarely observed in this study, as the clearance between

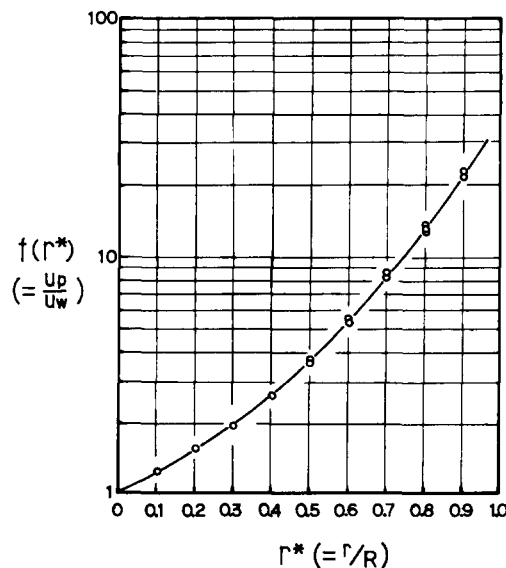


Figure 4. Experimental wall correction factors of single solid spheres in a triangular cylinder.

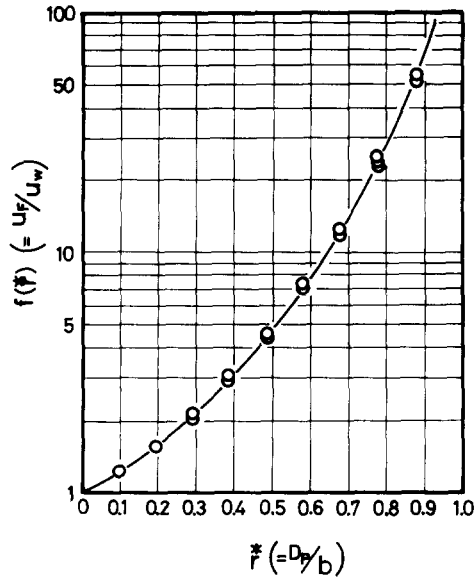


Figure 5. Wall correction factors of single solid spheres in a square cylinder.

the solid sphere of a fixed diameter and the inside wall of the triangular and a square cylinders and parallel plates is always larger than that between the solid sphere of the same diameter and the circular cylinder.

In the parallel plates the ball bearing sometimes settled with rotation along the centerline in the case of small clearances and these data were plotted by the triangular points in figure 6, but these phenomena were quite rare.

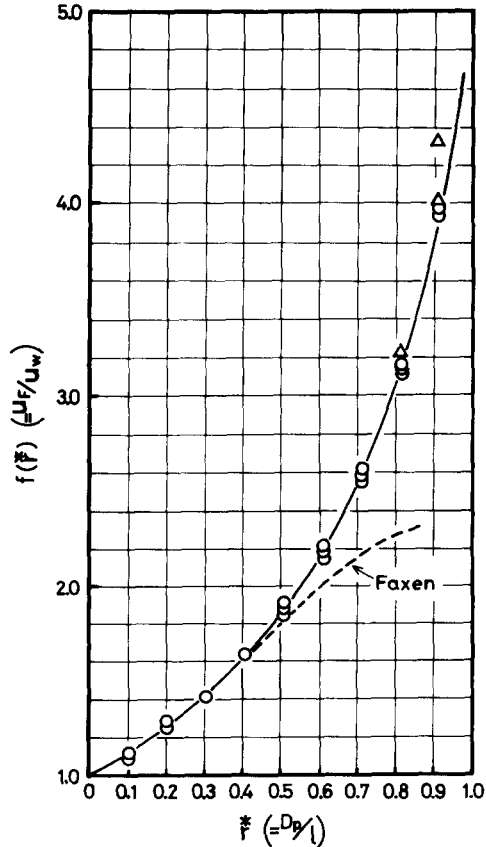


Figure 6. Wall correction factors of single solid spheres in parallel plates.

There is no study on the theoretical and experimental wall correction factors of the triangular and square cylinders, but Faxen (1922) obtained a theoretical wall correction factor for parallel plates by a series expansion solution, which is also plotted in Fig. 6. The series expansion solution by Faxen agreed very well with our experimental data up to $r^* = 0.4$, as it expanded a series in the vicinity of the center of $r^* = 0.0$. In this study the ratio of the diameter of the largest solid sphere to the longer width of the parallel plates is about 10 and hence the wall effect of the larger width is negligibly small.

The wall correction factors in the case of no rotation so obtained were correlated by the following 19th order equations by the least square method, the formulae of which are similar to the series expansion solution. The coefficients of the 19th order equations were tabulated table 1.

$$F(r)^* = \frac{1}{R_0 + R_1 r + R_2 r^2 + \dots + R_{17} r^{17} + R_{18} r^{18} + R_{19} r^{19}} \tag{2}$$

The standard variance is expressed by the following equation.

$$S = \sqrt{\left(\frac{1}{N} \sum_{i=1}^N [F_i(r)^* - F(r)^*]^2\right)}$$

The standard variances were 0.261, 0.485 and 0.0185 for the triangular and square cylinders and parallel plates, respectively and all the standard variances were considerably small, compared to the average wall correction factors, $F(r)$.

The triangular and square cylinders and parallel plates can be used as a falling ball viscometer from the wall correction factors so obtained. The viscometers of the triangular and square cylinders, and parallel plates can be easily made from the flat-transparent plates, which is more convenient than the conventional falling ball viscometer of the circular cylinder, of which the cross-section must be precisely circular. Furthermore, in the very small clearances

Table 1. Coefficients of polynomials by the least square method

	triangular cylinder	square cylinder	parallel plates
R ₀	1.0000000	1.0000000	1.0000000 10 ⁰
R ₁	-0.1524694 10 ¹	-0.1923777 10 ¹	-0.4027060 10 ⁰
R ₂	-0.9356945 10 ¹	0.1649393 10 ¹	-0.8435362 10 ¹
R ₃	0.6788950 10 ²	-0.1153624 10 ²	0.3487996 10 ²
R ₄	-0.1634936 10 ³	0.2682020 10 ²	-0.2359584 10 ²
R ₅	0.6563649 10 ²	0.1367386 10 ²	-0.1193919 10 ³
R ₆	0.1929998 10 ³	-0.5060226 10 ²	0.1362242 10 ³
R ₇	0.4729873 10 ²	-0.1042480 10 ³	0.1601959 10 ³
R ₈	-0.3751033 10 ³	0.1170802 10 ³	-0.4106427 10 ¹
R ₉	0.2752887 10 ³	0.2395431 10 ³	-0.3171554 10 ³
R ₁₀	-0.1190656 10 ⁴	-0.1757552 10 ³	-0.1989548 10 ³
R ₁₁	0.6542166 10 ³	0.1097079 10 ²	0.4608181 10 ²
R ₁₂	0.2075038 10 ⁴	-0.2409061 10 ³	0.3581750 10 ³
R ₁₃	0.3268518 10 ³	-0.9802373 10 ²	0.2128604 10 ³
R ₁₄	-0.2014247 10 ⁴	0.1344775 10 ³	0.2338137 10 ³
R ₁₅	-0.2869939 10 ⁴	0.2037087 10 ³	-0.8912624 10 ²
R ₁₆	0.1553641 10 ⁴	0.9298401 10 ²	-0.6472198 10 ²
R ₁₇	0.1483049 10 ⁴	-0.1153909 10 ³	-0.2528621 10 ³
R ₁₈	0.2078562 10 ⁴	0.2791606 10 ³	0.2000882 10 ³
R ₁₉	-0.2211987 10 ⁴	-0.3434466 10 ³	0.2897953 10 ³

the rotation phenomena rarely occurs in the triangular and square cylinders, and parallel plates, but always occur in the circular cylinder, which causes an inaccurate measurement of viscosity.

4. CONCLUSION

The wall correction factors of single solid spheres were experimentally measured in triangular and square cylinders, and parallel plates in the creeping flow region.

In the future the wall correction factors in such apparatus should be theoretically studied by series expansions and numerical methods and be compared with these experimental data.

NOMENCLATURE

- a the length of the perpendicular line from the center to a side of the triangular cylinder, half of the length of a side of the square cylinder or half of the shorter width of the parallel plates, cm
- $F(r^*)$ wall correction factor
- $F_i(r^*)$ measured wall correction factor
- N number of the data of the measured wall correction factors
- r radius of a solid sphere, cm
- r^* dimensionless radius of a solid sphere = r/a
- R_0-R_{19} coefficients of the 19th order equation
- S standard variance
- u_f free settling velocity of a single solid sphere calculated from Stokes law, cm/s
- u_w settling velocity of a single solid sphere influenced by wall, cm/s

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